

Dissipation mechanisms of the inner core's translational oscillations

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Probing of the Earth's interior is limited to a few measurements. For example, a more accurate estimate of the density difference between the inner and outer core could better constrain the driving force of the geodynamo. Gravimetric measurements of the **translational oscillations of the inner core** could help in this respect, but these oscillations still elude detection. **Translational oscillations**, also known as **Slichter modes**^{1,2}, are the result of extreme events, such as massive earthquakes or asteroid impacts, which can slightly **displace the inner core**. The centre of mass of the inner core would later swing around the equilibrium position as a **damped oscillator**.

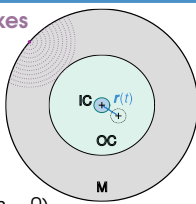
Previous linear models could only predict the **oscillation period**^{3,4}, bounding the frequency range of interest for observations. Here, for the first time, we study the **viscous and magnetic dissipation mechanisms** through non-linear simulations of the **outer core fluid response**. We take full advantage of the spherical shell geometry and use the fast pseudo-spectral code XSHELLS⁵ to solve the problem numerically. Since the study of realistic Earth values is out of reach, we use a systematic exploration of the parameter space to derive **scaling laws** that can be used to extrapolate to Earth conditions.

Translational oscillations^{1,2}

- Strong Earthquakes

- Impact events

can displace the inner core.



The **inner core** center of mass evolution $r(t)$ follows Newton's 2nd law^{3,4}:

$$m_{ic} \frac{d^2 r}{dt^2} = \sum_j F_j$$

- Coriolis
- centrifugal
- gravity
- added mass
- viscous
- magnetic

Modes splitting:

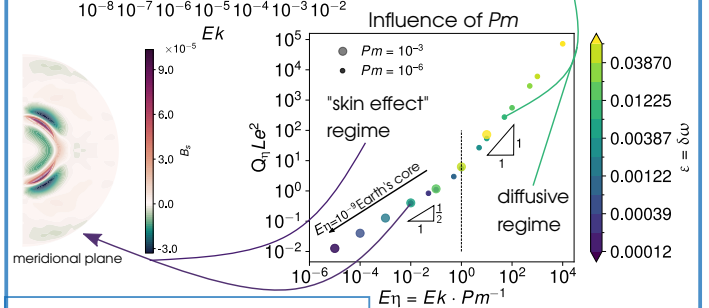
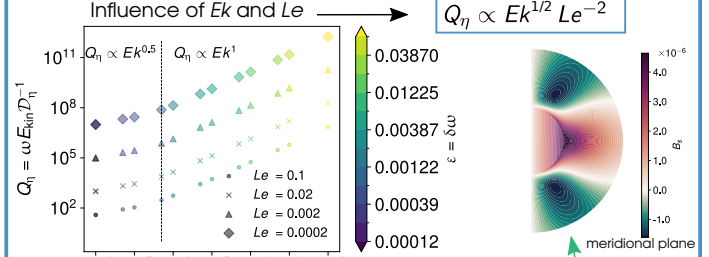
- 1 **polar** ($l = 1, m = 0$)
- 2 **equatorial** ($l = 1, m = -1, +1$)

Known **period** $T_{OSC} = 4-6$ h, but unclear **dissipation** mechanisms.

Magnetic dissipation of the polar mode

Forced **polar** oscillation: $u_z = \varepsilon \cos(\omega t)$

Uniform vertical magnetic field: $\mathbf{b} = B_0 \cdot \mathbf{e}_z$



our estimate for the Earth: $Q_\eta \approx 7 \cdot 10^3$
 Buffet et al. (1995)⁷: $2 \cdot 10^3 < Q_\eta < 5.8 \cdot 10^5$
Magnetic dissipation is the largest: $D_\eta \gg D_\nu$

Physical model implemented in XSHELLS⁵

Incompressible Navier-Stokes in the **outer core**

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} + 2\mathbf{e}_z \times \mathbf{u} = -\nabla p^* + Ek \nabla^2 \mathbf{u} + Le^2 (\nabla \times \mathbf{b}) \times \mathbf{b}, \quad \nabla \cdot \mathbf{u} = 0$$

Induction equation in the **inner** and the **outer core**

$$\frac{\partial \mathbf{b}}{\partial t} = \nabla \times (\mathbf{u} \times \mathbf{b}) + \frac{Ek Pm^{-1}}{E_\eta} \nabla^2 \mathbf{b}, \quad \nabla \cdot \mathbf{b} = 0$$

Velocity b.c. at the **inner core boundary** $u_r(R_i) = \varepsilon(t) Y_1^m(\theta, \varphi)$

Forced oscillations: $\varepsilon(t) = \varepsilon \cos(\omega t)$

OUTPUTS:

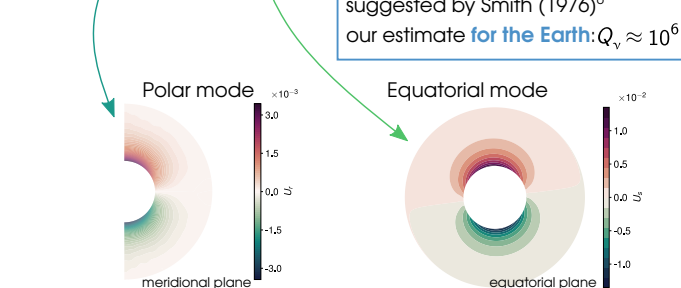
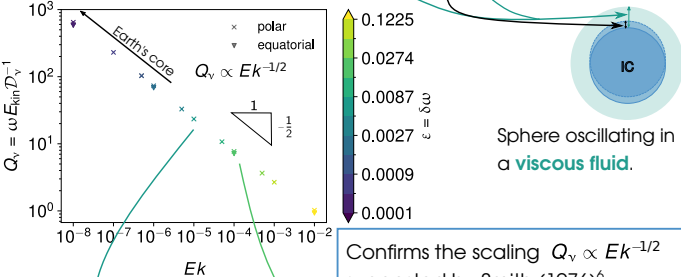
- kinetic energy: $E_k = \frac{1}{2} \int |\mathbf{u}|^2 dV$
- viscous dissipation: $D_\nu = \nu \int |\nabla \times \mathbf{u}|^2 dV$
- magnetic dissipation: $D_\eta = \eta_m \int |\nabla \times \mathbf{b}|^2 dV$
- quality factor: $Q_i = E_{kin} \omega D_i^{-1}$

INPUTS:

- dimensionless displacement: $\delta = \frac{\delta_{osc}}{R_o}$ **Earth core** 10^{-8}
- dimensionless frequency: $\omega = \frac{\omega_{osc}}{\Omega_e}$ **6**
- Ekman number: $Ek = \frac{\nu}{\Omega_e R_o^2}$ 10^{-15}
- magnetic Prandtl number: $Pm = \frac{\nu}{\eta_m}$ 10^{-6}
- Lehnert number: $Le = \frac{B_0}{\sqrt{\rho \mu_0 \Omega_e} R_o}$ 10^{-4}
- aspect ratio: $\Gamma = \frac{R_i}{R_o}$ **0.35**

Viscous dissipation

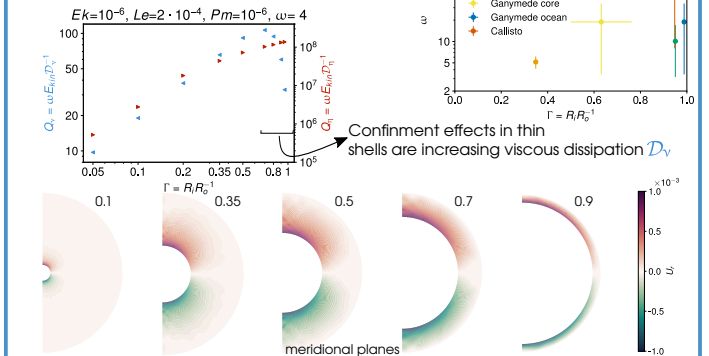
Oscillations inside the **viscous boundary layer**: $\delta < \delta_{\nu,bl} \propto \sqrt{Ek}$



Influence of the inner radius

Translational oscillations may occur in **planets**

and **moons**⁸ whenever a liquid shell is present.



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